

Turbulent transport via SPDEs and Wiener chaos

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The Transport Equation

$$\frac{\partial \theta(t, x)}{\partial t} = -\mathbf{v}(t, x) \cdot \nabla \theta(t, x), \quad t > 0, \quad x \in \mathbb{R}^d;$$

$$\theta(0, x) = \theta_0(x); \quad \mathbf{v} = (v^1, \dots, v^d) \in \mathbb{R}^d, \quad d \geq 2.$$

Laminar transport:

each v^i is Lipschitz continuous in x .

$$\theta(t, x) = \theta_0(X_{t,0}^x);$$

X is the flow of \mathbf{v} :

$$\frac{dX_{s,t}^x}{dt} = \mathbf{v}(t, X_{s,t}^x), \quad t > s, \quad X_{s,s}^x = x.$$

(Characteristic equation.)

Turbulent Transport

v is not Lipschitz continuous in x

- Example — Kolmogorov's theory: v is Hölder $\approx 1/3$ for $d = 3$.
- Difficulty — Existence but no uniqueness for the flow equation (intrinsic stochasticity).
- How to find θ ?

Regularization

- Introducing viscosity (κ -limit):

$$\frac{\partial \theta^\kappa(t, x)}{\partial t} = \kappa \Delta \theta^\kappa(t, x) - \mathbf{v}(t, x) \cdot \nabla \theta^\kappa(t, x), \quad t > 0, \quad x \in \mathbb{R}^d;$$

$$dX_{s,t}^{\kappa,x} = \mathbf{v}(t, X_{s,t}^{\kappa,x}) dt + \sqrt{2\kappa} dw(t), \quad t > s, \quad X_{s,s}^{\kappa,x} = x.$$

(Gawędzki and Vergassola (2000), Weinan E and Eric Vanden Eijnden (2000))

Unique strong solution of the flow equation!

- Smoothing out \mathbf{v} (ε -limit):

$$\mathbf{v}^\varepsilon(t, x) = \frac{1}{\varepsilon^d} \int_{\mathbb{R}^d} \mathbf{v}(t, y) \psi\left(\frac{x-y}{\varepsilon}\right) dy$$

(Weinan E and Eric Vanden Eijnden (2000))

Generalized Kraichnan's Model

Physical Model for \mathbf{v} :

- \mathbf{v} is a statistically homogeneous, isotropic, and stationary Gaussian vector field with zero mean and covariance

$$E(v^i(t, x)v^j(s, y)) = \delta(t - s)C^{ij}(x - y).$$

- For small x , $C^{ij}(x) \sim (\delta_{ij} - c^{ij}|x|^\gamma)$,
 $0 < \gamma < 2$.
- Div-free and rotation-free components of \mathbf{v} .
- Classification of flows in κ - and ε -limits.
Limits can be different!

Generalized Kraichnan's Model

Mathematical Model for \mathbf{v} .

(LeJan and Raimond (2002))

- The matrix C is characterized by its Fourier transform:

$$\widehat{C}(z) = \frac{A_0}{(1 + |z|^2)^{(d+\gamma)/2}} \left(a \frac{zz^*}{|z|^2} + \frac{b}{d-1} \left(I - \frac{zz^*}{|z|^2} \right) \right),$$

- $a = 0 \Rightarrow \nabla \cdot \mathbf{v} = 0$
(" b " — div-free component);
- $b = 0 \Rightarrow \mathbf{v} = \nabla V$ for some scalar V
(" a " — rotation-free component).
- H_C — Reproducing kernel Hilbert space for C .
 $H_C = H^{(d+\gamma)/2}(\mathbb{R}^d; \mathbb{R}^d)$ if $a, b > 0$.

- Representation of v :

$$v^i(t, x) = \sum_{k \geq 1} \sigma_k^i(x) \dot{w}_k(t);$$

$\dot{w}_k(t)$ are independent standard Gaussian white noises; $\{\sigma_k, k \geq 1\}$ is a CONS in H_C .

- σ_k^i is Hölder $\gamma/2$, $\sum_{k \geq 1} \sigma_k^i(x) \sigma_k^j(x) = \delta_{ij}$.

- Flow equation:

$$X^i(t; x, s) = x^i + \int_s^t \sigma_k^i(X(\tau; x, s)) dw_k(\tau).$$

- *Statistical solution* of the flow equation.
- $\zeta = b/(a + b)$ — degree of incompressibility.
 $\zeta = 0$ — Pure Gradient;
 $\zeta = 1$ — Incompressible.
- Classification of flows: phase diagram in (ζ, γ) plane.

Can be coalescing (particles stick together) or diffusive (particles split).

- Still very little info about θ .

Transport Equation as an SPDE

Notations: $D_i = \frac{\partial}{\partial x^i}$; $\Delta = D_i D_i$.

Summation convention: summation over a pair of repeating indices.

- $\mathbb{F} = (\Omega, \mathcal{F}, \{\mathcal{F}_t\}_{t \geq 0}, \mathbb{P})$, a stochastic basis with the usual assumptions.
- $(w_k(t), k \geq 1, t \geq 0)$, independent standard Wiener processes on \mathbb{F} .
- v divergence-free (incompressible flow).

Then

$$d\theta(t, x) = - \sum_k \sigma_k(x) \cdot \nabla \theta(t, x) \circ dw_k(t).$$

or

$$d\theta(t, x) = \frac{1}{2} \Delta \theta(t, x) dt - \sigma_k^i(x) D_i \theta(t, x) dw_k(t)$$

Fully Degenerate Equation:

$$\frac{1}{2} \sigma_k^i(x) \sigma_k^j(x) = \frac{1}{2} \delta_{ij}.$$

Wiener Chaos

$$W(t) = (w_k(t), k \geq 1, 0 \leq t \leq T).$$

$$\{m_i(s), i \geq 1\} \text{ — CONS in } L_2([0, T]),$$

$$\xi_{ik} = \int_0^T m_i(s) dw_k(s).$$

$$\mathcal{J} = \left\{ \alpha = (\alpha_i^k, i, k \geq 1) \mid |\alpha| = \sum_{i,k} \alpha_i^k < \infty \right\};$$

$$\xi_\alpha = \prod_{i,k} \left(\frac{H_{\alpha_i^k}(\xi_{ik})}{\sqrt{\alpha_i^k!}} \right), \text{ where}$$

$$H_n(x) = (-1)^n \exp\left\{\frac{x^2}{2}\right\} \frac{d^n}{dx^n} \exp\left\{-\frac{x^2}{2}\right\}.$$

$$\text{Example: } \alpha = \begin{pmatrix} 0 & 1 & 0 & 3 & 0 & 0 & \dots \\ 2 & 0 & 0 & 0 & 4 & 0 & \dots \\ 0 & 0 & 0 & 0 & 0 & 0 & \dots \\ \vdots & \vdots & \vdots & \vdots & \vdots & \vdots & \dots \end{pmatrix}$$

$$\alpha_2^1 = 1; \alpha_4^1 = 3; \alpha_1^2 = 2; \alpha_5^2 = 4.$$

$$\xi_\alpha = \xi_{2,1} \cdot \frac{H_3(\xi_{4,1})}{\sqrt{3!}} \cdot \frac{H_2(\xi_{1,2})}{\sqrt{2!}} \cdot \frac{H_4(\xi_{5,2})}{\sqrt{4!}}.$$

Theorem. (Cameron and Martin, 1947)

The collection $\{\xi_\alpha, \alpha \in \mathcal{J}\}$ is an orthonormal basis in $L_2(\Omega, \mathcal{F}_T^W, \mathbb{P})$:

If $\eta \in L_2(\Omega, \mathcal{F}_T^W, \mathbb{P})$ and $\eta_\alpha = \mathbb{E}(\eta\xi_\alpha)$, then

$$\eta = \sum_{\alpha \in \mathcal{J}} \eta_\alpha \xi_\alpha$$

and

$$\mathbb{E}|\eta|^2 = \sum_{\alpha \in \mathcal{J}} \eta_\alpha^2.$$

Goal: $\theta(t, x) = \sum_{\alpha \in \mathcal{J}} \theta_\alpha(t, x) \xi_\alpha$

The S-system

$$d\theta(t, x) = \frac{1}{2} \Delta \theta(t, x) dt - \sigma_k^i(x) D_i \theta(t, x) dw_k(t)$$

If σ_k^i, θ_0 are smooth, then $\theta(t, x) = \sum_{\alpha \in \mathcal{J}} \theta_\alpha(t, x) \xi_\alpha$.

Define: $\xi_\alpha(t) = \mathbb{E}(\xi_\alpha | \mathcal{F}_t^W)$; $\xi_\alpha(0) = I(|\alpha| = 0)$.

Fact: $d\xi_\alpha(t) = \sum_{i,k} \sqrt{\alpha_i^k} \xi_{\alpha^-(i,k)}(t) m_i(t) dw_k(t)$,

where $\alpha^-(i, k)$ is the multi-index with the components

$$\left(\alpha^-(i, k) \right)_j^l = \begin{cases} \max(\alpha_i^k - 1, 0), & \text{if } i = j \text{ and } k = l, \\ \alpha_j^l, & \text{otherwise.} \end{cases}$$

By the Itô formula

$$\begin{aligned} \frac{\partial \theta_\alpha(t, x)}{\partial t} &= \frac{1}{2} \Delta \theta_\alpha(t, x) \\ &\quad - \sum_{i,k} \sqrt{\alpha_i^k} \sigma_k^j(x) D_j \theta_{\alpha^-(i,k)}(t, x) m_i(t); \end{aligned}$$

$$\theta_\alpha(0, x) = \theta_0(x) I(|\alpha| = 0)$$

Solving the S-system

$|\alpha| = 0$:

$$\frac{\partial \theta_{(0)}(t, x)}{\partial t} = \frac{1}{2} \Delta \theta_{(0)}(t, x)$$

$$\theta_{(0)}(0, x) = \theta_0(x) \Rightarrow \theta_{(0)}(t, x) = \mathbb{T}_t \theta_0(x).$$

$\alpha = \delta_{ik}$:

$$\begin{aligned} \frac{\partial \theta_{(ik)}(t, x)}{\partial t} &= \frac{1}{2} \Delta \theta_{(ik)}(t, x) \\ &\quad - \sigma_k^j(x) D_j \theta_{(0)}(t, x) m_i(t); \quad \theta_{(ik)}(0, x) = 0. \end{aligned}$$

$$\theta_{(ik)}(t, x) = - \int_0^t m_i(s) \mathbb{T}_{t-s} \sigma_k^j D_j \mathbb{T}_s \theta_0(x) ds.$$

In fact, with $\mathcal{M}_k = -\sigma_k^j D_j$,

$$\begin{aligned} &\sum_{|\alpha|=N} |\theta_\alpha(t, x)|^2 \\ &= \sum_{k_1, \dots, k_N=1}^{\infty} \int_0^t \int_0^{s_N} \cdots \int_0^{s_2} \\ &\quad |\mathbb{T}_{t-s_N} \mathcal{M}_{k_N} \cdots \mathbb{T}_{s_2-s_1} \mathcal{M}_{k_1} \mathbb{T}_{s_1} \theta_0(x)|^2 ds_1 \cdots ds_N. \end{aligned}$$

Conclusion: σ_k^i do not have to be smooth or even continuous.

The Main Result

$$d\theta(t, x) = \frac{1}{2} \Delta \theta(t, x) dt - \sigma_k^i(x) D_i \theta(t, x) dw_k(t)$$

Theorem (Lototsky and Rozovskii, 2003)

If $\theta_0 \in L_2(\mathbb{R}^d)$, $D_i \sigma_k^i = 0$, and $\sigma_k^i(x) \sigma_k^j(x) = \delta_{ij}$, then

- For $t > 0$, $\sum_{\alpha \in \mathcal{J}} \|\theta_\alpha(t)\|_{L_2(\mathbb{R}^d)}^2 \leq \|u_0\|_{L_2(\mathbb{R}^d)}^2$.
- For every $\varphi \in C_0^\infty(\mathbb{R}^d)$, the random field $\theta(t, x) = \sum_{\alpha \in \mathcal{J}} \theta_\alpha(t, x) \xi_\alpha$ satisfies

$$\begin{aligned} (\theta, \varphi)(t) &= (\theta_0, \varphi) + \frac{1}{2} \int_0^t (\theta, \Delta \varphi)(s) ds \\ &\quad + \int_0^t (\theta, \sigma_k^i D_i \varphi) dw_k(s). \end{aligned}$$

- If

$$X^i(s; t, x) = x^i - \int_s^t \sigma_k^i(X(s; t, x)) \overleftarrow{dw}_k(s), \quad s \in [0, t],$$

(weak solution), then

$$\theta(t, x) = \mathbb{E} \left(\theta_0(X(0; t, x)) \mid \mathcal{F}_t^W \right)$$

Discussion

- Transport Equation \Rightarrow The S-system \Rightarrow Coefficients $\theta_\alpha(t, x) \Rightarrow$

$$\text{Solution } \theta(t, x) = \sum_{\alpha \in \mathcal{J}} \theta_\alpha(t, x) \xi_\alpha.$$

- Probabilistically strong solution.
- $\mathbb{E}\theta(t, x) = \theta_{(0)}(t, x).$
- $\mathbb{E}(\theta(t, x)\theta(s, y)) = \sum_{\alpha \in \mathcal{J}} \theta_\alpha(t, x)\theta_\alpha(s, y).$
- By interpolation: $\mathbb{E}\|\theta\|_{L_p(\mathbb{R}^d)}^p \leq \|\theta_0\|_{L_p(\mathbb{R}^d)}^p$, $2 < p < \infty$. Weighted L_p (e.g. $\theta_0(x) = |x|$) are also OK.
- Incompressible flow \Rightarrow Itô = Stratonovich.
- Conservation of energy \Leftrightarrow Strong solution of the flow equation.

Totally Turbulent Transport

$$d\theta(t, x) = \nu \Delta \theta(t, x) dt - \sigma_k^i(x) D_i \theta(t, x) dw_k(t),$$

σ_k , $k \geq 1$ — CONS in $L_2(\mathbb{R}^d; \mathbb{R}^d)$.
(P. Chow, J. Potthoff, etc.)

Note: $\sum_{k \geq 1} \sigma_k^i(x) \sigma_k^j(x)$ diverges.

S-system:

$$\begin{aligned} \frac{\partial \theta_\alpha(t, x)}{\partial t} &= \nu \Delta \theta_\alpha(t, x) \\ &\quad - \sum_{i,k} \sqrt{\alpha_i^k} \sigma_k^j(x) D_j \theta_{\alpha-(i,k)}(t, x) m_i(t); \end{aligned}$$

Still solvable, but now

$$\sum_{\alpha \in \mathcal{J}} \|\theta_\alpha(t)\|_{L_2(\mathbb{R}^d)}^2 = +\infty.$$

Weighted Wiener Chaos

Idea: replace w_k with $q_k w_k$.

$Q = \{q_1, q_2, \dots\}$, $q_k > 0$; $Q \leq \tilde{Q}$ means $q_k \leq \tilde{q}_k$ for all k ;

$$q^\alpha = \prod_{i,k} q_k^{\alpha_i^k}.$$

Definition. For a Banach space X , the Q -weighted Wiener Chaos space $L_{2,Q}(\mathcal{F}_T^W; X)$ is

$$L_{2,Q}(\mathcal{F}_T^W; X) = \left\{ (u_\alpha) : \sum_{\alpha \in \mathcal{J}} q^{2\alpha} \|u_\alpha\|_X^2 < \infty \right\}.$$

Still write $u = \sum_{\alpha \in \mathcal{J}} u_\alpha \xi_\alpha$

$$Q = 1 \Rightarrow L_{2,Q}(\mathcal{F}_T^W; X) = L_2(\mathcal{F}_T^W; X);$$

$$Q \leq \tilde{Q} \Rightarrow L_{2,\tilde{Q}}(\mathcal{F}_T^W; X) \subseteq L_{2,Q}(\mathcal{F}_T^W; X).$$

Q-transform:

$$u = \sum_{\alpha \in \mathcal{J}} u_\alpha \xi_\alpha \leftrightarrow u^Q = \sum_{\alpha \in \mathcal{J}} q^\alpha u_\alpha \xi_\alpha$$

$$du = A u dt + \sum_{k \geq 1} \mathcal{B}_k u dw_k$$

$$\Updownarrow$$

$$du^Q = A u^Q dt + \sum_{k \geq 1} q_k \mathcal{B}_k u^Q dw_k$$

Example: $w_k \rightarrow q_k w_k$

1. (Obvious) If $u(t) = 1 + \sum_{k \geq 1} \int_0^t u(s) dw_k(s)$,

then $u \in L_{2,Q}(\mathcal{F}_T^W; \mathbb{R})$ for every $Q = (q_1, q_2, \dots)$
so that

$$\sum_{k \geq 1} q_k^2 < \infty.$$

2. (Nualart-Rozovskii, 1997) If

$$du(t, x) = \Delta u(t, x) dt + u(t, x) dw(t, x),$$

$$t > 0, x \in \mathbb{R}^d, d \geq 2,$$

then $u \in L_{2,Q}(\mathcal{F}_T^W; L_2(\mathbb{R}^d))$ for some Q .

Theorem (Lototsky and Rozovskii, 2004)

Assume that $\theta_0 \in L_2(\mathbb{R}^d)$ and $|\sigma_k^i(x)| \leq C_k$.

Let Q be a sequence with $q_k = \frac{\sqrt{\delta\nu}}{d2^k C_k}$ for some $0 < \delta < 2$. If

$$\begin{aligned} \frac{\partial \theta_\alpha(t, x)}{\partial t} &= \nu \Delta \theta_\alpha(t, x) \\ &\quad - \sum_{i,k} \sqrt{\alpha_i^k} \sigma_k^j(x) D_j \theta_{\alpha-(i,k)}(t, x) m_i(t); \end{aligned}$$

then $\sum_{\alpha \in \mathcal{J}} q^{2\alpha} \|\theta_\alpha(t)\|_{L_2(\mathbb{R}^d)}^2 < \infty$

and $\theta(t, x) = \sum_{\alpha \in \mathcal{J}} \theta_\alpha(t, x) \xi_\alpha$ satisfies

$$\theta \in L_{2,Q} \left(\mathcal{F}_T^W; \mathbf{C}((0, T); L_2(\mathbb{R}^d)) \right).$$

This θ is called the *Wiener Chaos solution* of the totally turbulent transport equation

$$d\theta(t, x) = \nu \Delta \theta(t, x) dt - \sigma_k^i(x) D_i \theta(t, x) dw_k(t).$$

Wiener Chaos Approach

- Equation for $u = u(t) \Rightarrow$ The S-system \Rightarrow Coefficients $u_\alpha(t) \Rightarrow$

$$\text{Solution } u(t) = \sum_{\alpha} u_{\alpha}(t) \xi_{\alpha}.$$

- Probabilistically strong solution.
- Computable expressions for the solution and its moments from the S-system.
- New regularity results.
- Possibilities for generalization.

Further Directions

- Anticipating equations.
- Elliptic equations.
- Nonlinear equations.